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The effects of deformation parameter on thermal width of moving quarkonia in plasma

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ABSTRACT: In general we can say that the thermal width of quarkonia corresponds to imaginary part of it's potential. Gravity dual of theories give explicit form of potential as $V_{Q\bar{Q}}$. Since there is an explicit formula for $ImV_{Q\bar{Q}}$ one can consider different gravity duals and study the results of contribution of various parameters. Variable gravity duals of moving pair in plasma have different results for potential. Our paper shows that deformation parameter c in warp factor leads to new results that we present them for arbitrary angles of the pair with respect to it's velocity. We compare our results with the case that no deformation parameter is in metric background. We will see that the thermal width of the pair increases with increasing deformation parameter. Also, for nonzero values of deformation parameter the pair feels moving plasma in all distances. In addition our results indicate that contribution of deformation parameter leads to larger dissociation length for the moving pair reverse to the effect of the pair's velocity in the plasma.

KEYWORDS: Gauge-gravity correspondence, AdS-CFT Correspondence, Holography and quark-gluon plasmas

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Contents

1	Introduction	1
2	$Q\bar{Q}$ in an deformed AdS, perpendicular case	2
3	$Q\bar{Q}$ in an deformed AdS at arbitrary angles	7
4	Concolusion	14

1 Introduction

When we want to study $Q\bar{Q}$ interaction we should consider the effect of the medium in motion of $Q\bar{Q}$, because this pair is not produced at rest in QGP. So, the velocity of the pair through the plasma has some effects on its interactions that should be taken into account. The interaction energy has a finite imaginary part at finite temperature that can be used to estimate the thermal width of the quarkonia [1, 2]. Calculations of $\text{Im}V_{Q\bar{Q}}$ relevant to QCD and heavy ion collisions were performed for static $Q\bar{Q}$ pairs using pQCD [3] and lattice QCD [4–6] before AdS/CFT.

The AdS/CFT is a correspondence [7–10] between a string theory in AdS space and a conformal field theory in physical space-time. It leads to an analytic semi-classical model for strongly coupled QCD. It has scale invariance, dimensional counting at short distances and color confinement at large distances. This theory describes the phenomenology of hadronic properties and demonstrate their ability to incorporate such essential properties of QCD as confinement and chiral symmetry breaking. In the AdS/CFT point of view the AdS_5 plays important role in describing QCD phenomena. So in order to describe a confining theory, the conformal invariance of AdS_5 must be broken somehow. Two strategies AdS/QCD background have been suggested in the literatures hard-wall model [11–16] and soft-wall model [17–35]. In hard-wall model to impose confinement and discrete normalizable modes that is to truncate the regime where string modes can propagate by introducing an IR cutoff in the fifth dimension at a finite value $z_0 \sim \frac{1}{\Lambda_{QCD}}$. Thus, the hard-wall at z_0 breaks conformal invariance and allows the introduction of the QCD scale and a spectrum of particle states, they have phenomenological problems, since the obtained spectra does not have Regge behavior. To remedy this it is necessary to introduce a soft cut off, using a dilaton field or using a warp factor in the metric [11, 22]. These models are called soft wall models. The soft-wall and hard-wall approach has been successfully applied to the description of the mass spectrum of mesons and baryons, the pion leptonic constant, the electromagnetic form factors of pion and nucleons, etc. On the other hand the study of the moving heavy quarkonia in space-time with AdS/QCD approach plays important role in interaction energy [36–39]. By using different metric backgrounds we see different effects on interaction energy.

Evaluation of $\text{Im}V_{Q\bar{Q}}$ will yield to determine the suppression of $Q\bar{Q}$ in heavy ion collision [40]. The main idea is using boosted frame to have $\text{Re} V_{Q\bar{Q}}$ and $\text{Im} V_{Q\bar{Q}}$ [41] for $Q\bar{Q}$ in a plasma.

From viewpoint of holography, the AdS/CFT correspondence can describe a “broken conformal symmetry”, when one adds a proper deformed warp factor in front of the AdS_5 metric structure [17, 42–61]. So, e^{cz^2} is a positive quadratic correction with z , the fifth dimension.

One natural question is about the connection between the warp factor and the potential $V_{Q\bar{Q}}$. In this work, the procedure of [40] is followed to evaluate the imaginary part of potential for an AdS metric background with deformation parameter in warp factor. It is interesting to see “what will happen if meson be in a deformed AdS?”

It is a trend to see the effects of deformation parameter on $\text{Re}V_{Q\bar{Q}}$ and $\text{Im}V_{Q\bar{Q}}$ which are evidences for “usual” or “unusual” behavior of meson in compare with the $c = 0$ case. As expected in the limit of $c \rightarrow 0$, all results are equal to the results of AdS_5 case. All above informations give us motivation to work on effect of the deformation parameter in AdS metric background on real and imaginary parts of potential. So, we organized the paper as follows. In section 2, we discuss the case where the pair is moving perpendicularly to the joining axis of the dipole in deformed AdS, we assume this metric background for $Q\bar{Q}$ and find some relations for real and imaginary parts of potential. This example will be presented with some numerical results for different values of deformation parameter. Then we consider general orientation of $Q\bar{Q}$ in section 3 and follow the procedure as before. Section 4 would be our conclusion and some suggestions for future work.

2 $Q\bar{Q}$ in an deformed AdS, perpendicular case

In this section we consider soft-wall metric background with deformation parameter in warp factor at finite temperature case. So, we present general relations for real and imaginary parts of potential when the dipole is moving with velocity η perpendicularly to the wind [40].

In our case we apply the general result for deformed AdS, the dual gravity metric will be as:

$$ds^2 = e^{2A(z)} \left[-f(z)dt^2 + \sum_{i=1}^3 dx^2 + \frac{1}{f(z)} dz^2 \right], \quad (2.1)$$

Where $A(z) = -\ln \frac{z}{R} + \frac{1}{4}cz^2$ and $f(z) = 1 - (\frac{z}{z_h})^4$. As mentioned before c is deformation parameter and R is the AdS curvature radius, also $0 \leq z \leq z_h$, $z_h = \frac{1}{\pi T}$ and T is boundary field theory’s temperature. We have a dynamic dilaton in action for the background and we write our calculations in string frame. If dilaton is such that it enters directly in the worldsheet action in the form ϕR . May be our concern is about the effect of a nontrivial dilaton profile to the string action. But somewhen people neglect it at the first step [62] and leave it for future study. Then one can check that the integral on the action correspond to worldsheet with higher genus. This means that we are doing string interactions and going to higher order in string perturbation theory. But now, for leading order calculations in genus, we need not to bother with this term even if the geometry has a dynamical dilaton. On the other hand one trace of dynamical dilaton can appear via temperature if we want

to calculate it with [63] approach. So, the exact temperature will be in hand. But we refer the reader to [61] for the reasons that in deformed AdS model with quadratic correction in warp factor the “temperature” takes the form of AdS-SW BH temperature. So, we have a deformed AdS which in the limit $c \rightarrow 0$ becomes AdS_5 . This comparing results help us to understand the effects of deformation parameter on the physical quantities such as interaction energy. Our calculations in the cases of LT , $ReV_{Q\bar{Q}}$ and $ImV_{Q\bar{Q}}$ give us motivation to compare results between different values of deformation parameter.

From metric background (2.1) one can obtain:

$$G_{00} = \frac{R^2}{z^2} \left[1 - \left(\frac{z}{z_h} \right)^4 \right] e^{\frac{cz^2}{2}} \quad (2.2)$$

$$G_{xx} = \frac{R^2}{z^2} e^{\frac{cz^2}{2}} \quad (2.3)$$

$$G_{zz} = \frac{R^2}{z^2} \left[1 - \left(\frac{z}{z_h} \right)^4 \right]^{-1} e^{\frac{cz^2}{2}}, \quad (2.4)$$

with these definitions,

$$\tilde{M}(z) \equiv M(z) \cosh^2 \eta - N(z) \sinh^2 \eta \quad (2.5)$$

$$\tilde{V}(z) \equiv V(z) \cosh^2 \eta - P(z) \sinh^2 \eta \quad (2.6)$$

$$M(z) \equiv G_{00} G_{zz} \quad (2.7)$$

$$V(z) \equiv G_{00} G_{xx} \quad (2.8)$$

$$P(z) \equiv G_{xx}^2 \quad (2.9)$$

$$N(z) \equiv G_{xx} G_{zz}, \quad (2.10)$$

we continue with hamiltonian,

$$H(z) \equiv \sqrt{\frac{\tilde{V}(z)}{\tilde{V}_*} \frac{\tilde{V}(z) - \tilde{V}_*}{\tilde{M}(z)}}, \quad (2.11)$$

where \tilde{V}_* means $\tilde{V}(z_*)$ and $*$ is the deepest position of the string in the bulk.

The equation of motion and the boundary conditions of the string relates L (length of the line joining both quarks) with z_* as follows,

$$\frac{L}{2} = \int_{r_*}^{\Lambda} \frac{dr}{H(r)}. \quad (2.12)$$

So, for the corresponding case we have,

$$\frac{L}{2} = - \int_0^{z_*} \frac{dz}{H(z)}. \quad (2.13)$$

In order to relation between S_{str} and z_* we find the regularized integral [41] as,

$$\begin{aligned} S_{\text{str}}^{\text{reg}} = & \frac{T}{\pi\alpha'} \int_{r_*}^{\infty} dr \left[\sqrt{\tilde{M}(r)} \sqrt{\frac{\tilde{V}(r)}{\tilde{V}(r_*)}} \left(\frac{\tilde{V}(r)}{\tilde{V}(r_*)} - 1 \right)^{-1/2} - \sqrt{M_0(r)} \right] \\ & - \frac{T}{\pi\alpha'} \int_{r_h}^{r_*} dr \sqrt{M_0(r)}, \end{aligned} \quad (2.14)$$

and we obtain the following results

$$LT = \frac{2}{\pi} y_h \sqrt{1 - y_h^4 \cosh^2 \eta} \int_1^\infty \frac{dy}{\sqrt{(y^4 - y_h^4) \left[e^{\frac{cy_h^2}{\pi^2 T^2} \left(\frac{1}{y^2} - 1 \right)} (y^4 - y_h^4 \cosh^2 \eta) - (1 - y_h^4 \cosh^2 \eta) \right]}} \quad (2.15)$$

where $y = \frac{z_*}{z}$ and $y_h = \frac{z_*}{z_h}$

$$S_{\text{str}}^{\text{reg}} = T^2 \frac{\sqrt{\lambda}}{y_h} \left\{ \int_1^\infty dy \left[\frac{e^{\frac{cy_h^2}{\pi^2 T^2} \left(\frac{1}{y^2} - \frac{1}{2} \right)} (y^4 - y_h^4 \cosh^2 \eta)}{\sqrt{(y^4 - y_h^4) \left[e^{\frac{cy_h^2}{\pi^2 T^2} \left(\frac{1}{y^2} - 1 \right)} (y^4 - y_h^4 \cosh^2 \eta) - (1 - y_h^4 \cosh^2 \eta) \right]}} \right] - e^{\frac{cy_h^2}{\pi^2 T^2 y^2}} - \int_0^1 dy e^{\frac{cy_h^2}{\pi^2 T^2 y^2}} \right\}, \quad (2.16)$$

Where $\lambda = \frac{R^4}{\alpha'^2}$ is and α' is the 't Hooft coupling of the gauge theory. Finally, we find the real part of potential as

$$ReV_{Q\bar{Q}} = \frac{S_{\text{str}}^{\text{reg}}}{T}.$$

Now we present a derivation of relation for imaginary part of potential from [41]. The reader can see more details in that reference. From there we can say one should consider the effect of worldsheet fluctuations around the classical configuration $z_c(x)$,

$$z(x) = z_c(x) \rightarrow z(x) = z_c(x) + \delta z(x). \quad (2.17)$$

And then the fluctuations should be taken into account in partition function so one arrives at,

$$Z_{\text{str}} \sim \int \mathcal{D}\delta z(x) e^{iS_{NG}(z_c(x) + \delta z(x))}. \quad (2.18)$$

Then there is an imaginary part of potential in action so, by dividing the interval region of x into $2N$ points where $N \rightarrow \infty$ that should be taken into account at the end of calculation we arrive at,

$$Z_{\text{str}} \sim \lim_{N \rightarrow \infty} \int d[\delta z(x_{-N})] \dots d[\delta z(x_N)] \exp \left[i \frac{\mathcal{T} \Delta x}{2\pi\alpha'} \sum_j \sqrt{M(z_j)(z'_j)^2 + V(z_j)} \right]. \quad (2.19)$$

Notice that we should expand $z_c(x_j)$ around $x = 0$ and keep only terms up to second order of it because thermal fluctuations are important around z_* which means $x = 0$,

$$z_c(x_j) \approx z_* + \frac{x_j^2}{2} z_c''(0), \quad (2.20)$$

With considering small fluctuations finally we will have,

$$V(z_j) \approx V_* + \delta z V'_* + z_c''(0) V'_* \frac{x_j^2}{2} + \frac{\delta z^2}{2} V''_*, \quad (2.21)$$

where $V_* \equiv V(z_*)$ and $V'_* \equiv V'(z_*)$. With (2.20), (2.21) and (2.19) one can derive (2.22), (2.23) and (2.24),

$$S_j^{NG} = \frac{\mathcal{T} \Delta x}{2\pi\alpha'} \sqrt{C_1 x_j^2 + C_2} \quad (2.22)$$

$$C_1 = \frac{z_c''(0)}{2} [2M_* z_c''(0) + V'_*] \quad (2.23)$$

$$C_2 = V_* + \delta z V'_* + \frac{\delta z^2}{2} V''_*. \quad (2.24)$$

For having $\text{Im} V_{Q\bar{Q}} \neq 0$ the function in the square root of (2.22) should be negative. then, we consider j-th contribution to Z_{str} as,

$$I_j \equiv \int_{\delta z_{j\min}}^{\delta z_{j\max}} d(\delta z_j) \exp \left[i \frac{\mathcal{T} \Delta x}{2\pi\alpha'} \sqrt{C_1 x_j^2 + C_2} \right], \quad (2.25)$$

For every δz between minimum and maximum of it's values which are the roots of $C_1 x_j^2 + C_2$ in δz , one leads to $C_1 x_j^2 + C_2 < 0$. The external value of the function

$$D(\delta z_j) \equiv C_1 x_j^2 + C_2(\delta z_j) \quad (2.26)$$

is,

$$\delta z = -\frac{V'_*}{V''_*}. \quad (2.27)$$

So, $D(\delta z_j) < 0 \longrightarrow -x_c < x_j < x_c$ leads us to have an imaginary part in square root, where,

$$x_c = \sqrt{\frac{1}{C_1} \left[\frac{V'^2_*}{2V''_*} - V_* \right]}. \quad (2.28)$$

If the square root in (2.28) is not real we should take $x_c = 0$. After all these conditions we can approximate $D(\delta z)$ by $D(-\frac{V'_*}{V''_*})$ in I_j ,

$$I_j \sim \exp \left[i \frac{\mathcal{T} \Delta x}{2\pi\alpha'} \sqrt{C_1 x_j^2 + V_* - \frac{V'^2_*}{2V''_*}} \right]. \quad (2.29)$$

The total contribution to the imaginary part, will be in hand with continuum limit. So,

$$\text{Im} V_{Q\bar{Q}} = -\frac{1}{2\pi\alpha'} \int_{|x| < x_c} dx \sqrt{-x^2 C_1 - V_* + \frac{V'^2_*}{2V''_*}}. \quad (2.30)$$

And finally after evaluating the integral one can arrive at the expression for imaginary part of potential as,

$$\text{Im} V_{Q\bar{Q}} = -\frac{1}{2\sqrt{2}\alpha'} \sqrt{M_*} \left[\frac{V'_*}{2V''_*} - \frac{V_*}{V'_*} \right]. \quad (2.31)$$

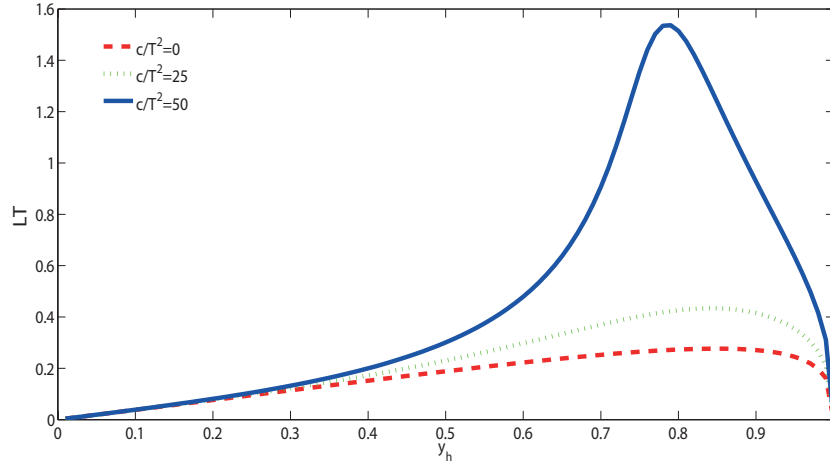


Figure 1. LT as a function of y_h at $\eta = 0$. $Q\bar{Q}$ is oriented to the hot wind and different values of deformation parameter are contributed. The solid blue curve corresponds to $\frac{c}{T^2} = 50$, the dotted green curve to $\frac{c}{T^2} = 25$ and the dashed red curve to $\frac{c}{T^2} = 0$.

Now, we are ready to calculate the imaginary part of potential in case of $\tilde{M}_* > 0$ so according to (2.31) and with our deformed AdS metric (2.1) we have following relation,

$$\begin{aligned} \frac{ImV_{Q\bar{Q}}}{\sqrt{\lambda}} = & -\frac{\pi T}{4\sqrt{2}y_h} e^{\frac{cy_h^2}{2\pi^2 T^2}} \sqrt{\frac{1 - y_h^4 \cosh^2 \eta}{1 - y_h^4}} \\ & \times \left[\frac{\frac{2cy_h^2}{\pi^2 T^2} (1 - y_h^4 \cosh^2 \eta) + 4}{\frac{cy_h^2}{\pi^2 T^2} (2 - 10y_h^4 \cosh^2 \eta) + 12} - \frac{(1 - y_h^4 \cosh^2 \eta)}{2 + \frac{cy_h^2}{2\pi^2 T^2} (1 - y_h^4 \cosh^2 \eta)} \right]. \end{aligned} \quad (2.32)$$

In figures 1 and 2 we can see the behavior of LT as a function of y_h for different values of deformation parameter for this perpendicular case. As we show the maximum of the $LT(y_h)$ which is an indicative of the limit of classical gravity calculations, increases with increasing deformation parameter.

On the contrary, increasing velocity reduces LT_{\max} , as it is mentioned in figure 3. Furthermore, increasing deformation parameter increases LT_{\max} which has been used to define a dissociation length for the moving $Q\bar{Q}$ pair.

Figures 4, 5 and 6 show the behavior of $ReV_{Q\bar{Q}}$ as a function of LT . As we know for zero value of c in short distances the pair does not feel the moving plasma and upper branch shows saddle point of string action. In compare with that, when c is nonzero pair can feel moving plasma for all values of distances. In addition, we can see with increasing deformation parameter the real part of potential increases and the unphysical curve corresponds to $q < q_{\max}$ and the lower branch is the dominant contribution to the action which corresponds to $q > q_{\max}$.

The imaginary part of the potential corresponds to the dissociation properties of heavy quarkonia. In figure 7 our results indicate that the thermal width of the pair increases with increasing the deformation parameter at a fixed velocity.

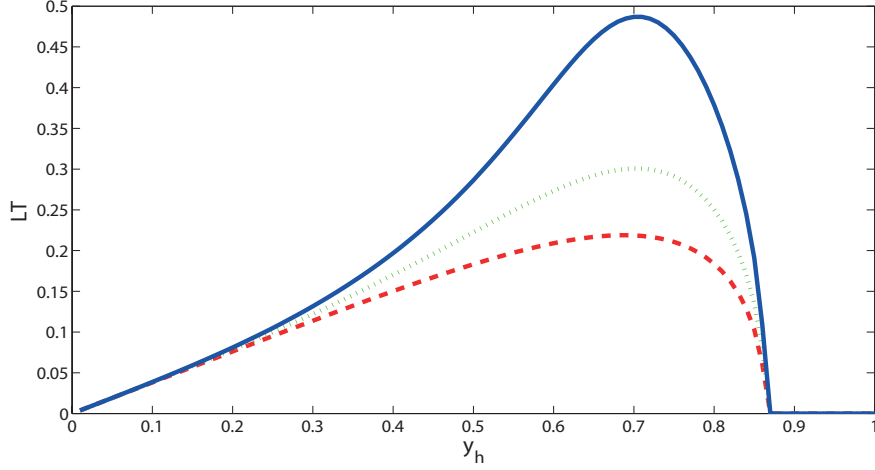


Figure 2. LT as a function of y_h at $\eta = 0.8$. $Q\bar{Q}$ is oriented to the hot wind and different values of deformation parameter are contributed. The solid blue curve corresponds to $\frac{c}{T^2} = 50$, the dotted green curve to $\frac{c}{T^2} = 25$ and the dashed red curve to $\frac{c}{T^2} = 0$.

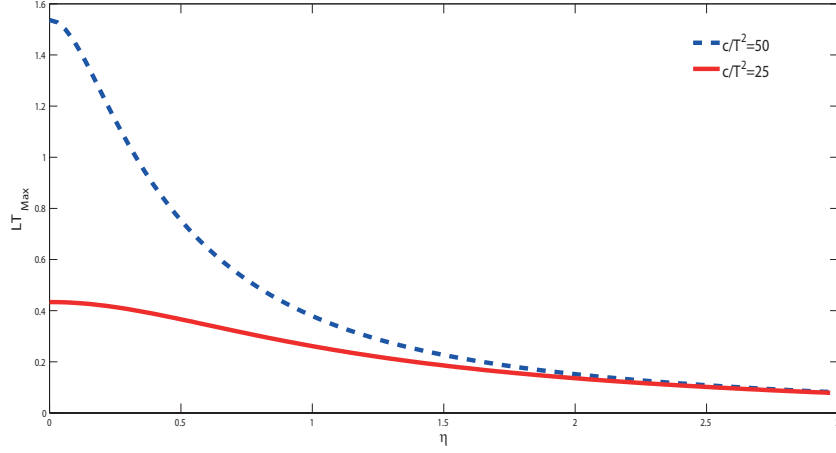


Figure 3. LT_{\max} as a function of η . $Q\bar{Q}$ is oriented to the hot wind and different values of deformation parameter are contributed. The dashed blue curve corresponds to $\frac{c}{T^2} = 50$ and the solid red curve to $\frac{c}{T^2} = 25$.

3 $Q\bar{Q}$ in an deformed AdS at arbitrary angles

In this section we extend our calculations for arbitrary angles, it means that orientation of dipole can have any arbitrary angle with respect to velocity vector. As before we extract the real and imaginary parts of potential with method of [40], θ is the angle of the dipole with respect to the X_{d-1} and dipole is on the (X_1, X_{d-1}) plane. The boundary conditions are,

$$\begin{aligned} z \left(\pm \frac{L}{2} \sin \theta \right) &= 0 \\ X_d \left(\pm \frac{L}{2} \sin \theta \right) &= \pm \frac{L}{2} \cos \theta \end{aligned} \quad (3.1)$$

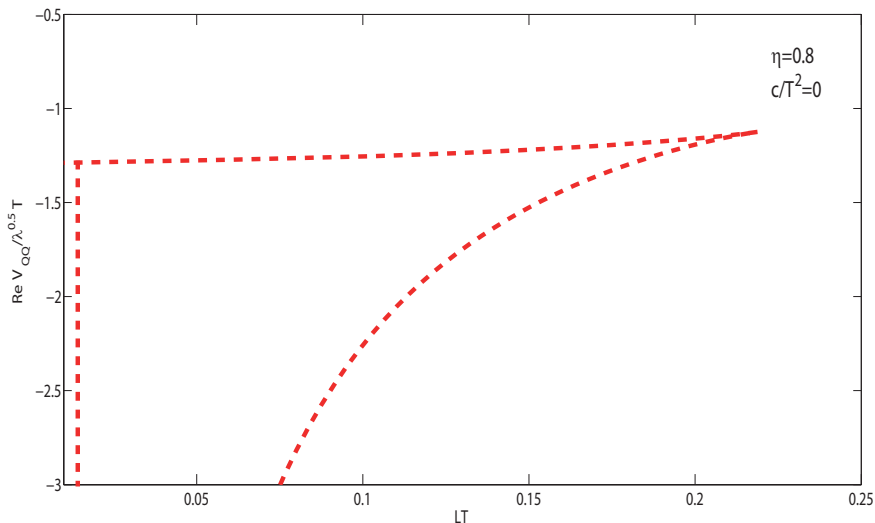


Figure 4. $ReV_{Q\bar{Q}}$ as a function of LT at a fixed velocity $\eta = 0.8$ the pair is oriented to the hot wind and deformation parameter is zero.

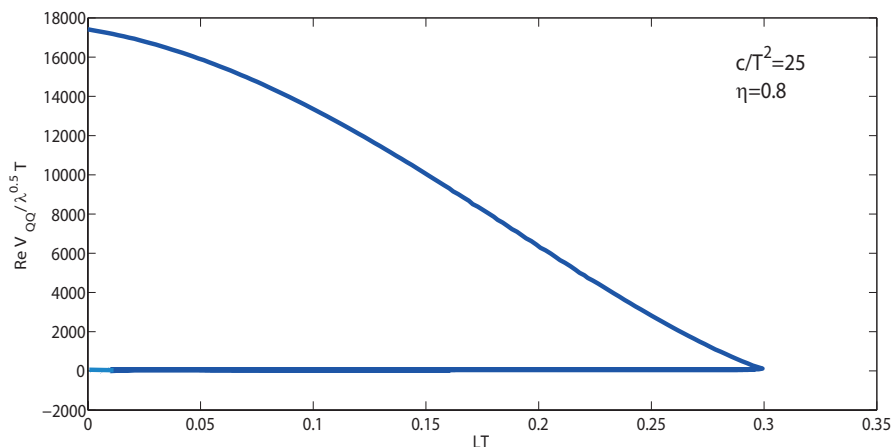


Figure 5. $ReV_{Q\bar{Q}}$ as a function of LT at a fixed velocity $\eta = 0.8$ the pair is oriented to the hot wind and scaled deformation parameter is 25.

And the action is,

$$S_{\text{str}} = -\frac{\mathcal{T}}{2\pi\alpha'} \int d\sigma \mathcal{L}, \quad (3.2)$$

where the lagrangian is defined as,

$$\mathcal{L} \equiv \sqrt{[M(z) \cosh^2 \eta - N(z) \sinh^2 \eta] z'(\sigma)^2 + V(z) X_d'(\sigma)^2 + [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta]}. \quad (3.3)$$

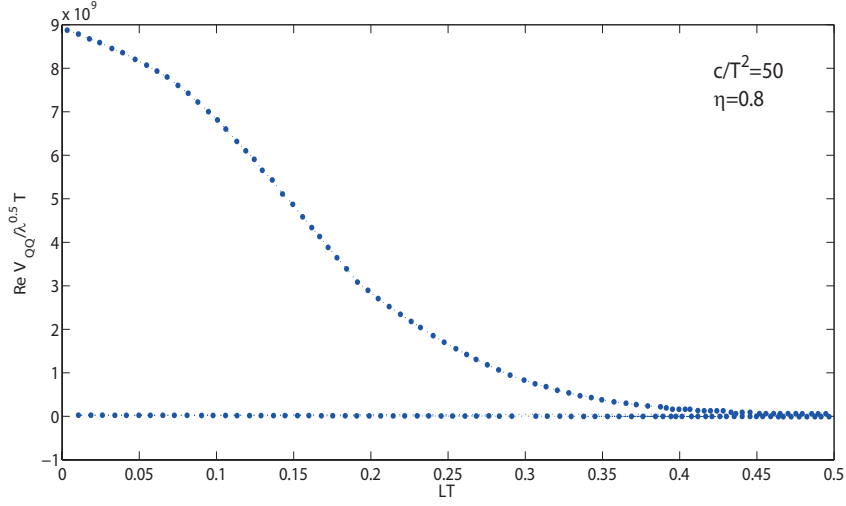


Figure 6. $\text{Re} V_{Q\bar{Q}}$ as a function of LT at a fixed velocity $\eta = 0.8$ the pair is oriented to the hot wind and scaled deformation parameter is 50.

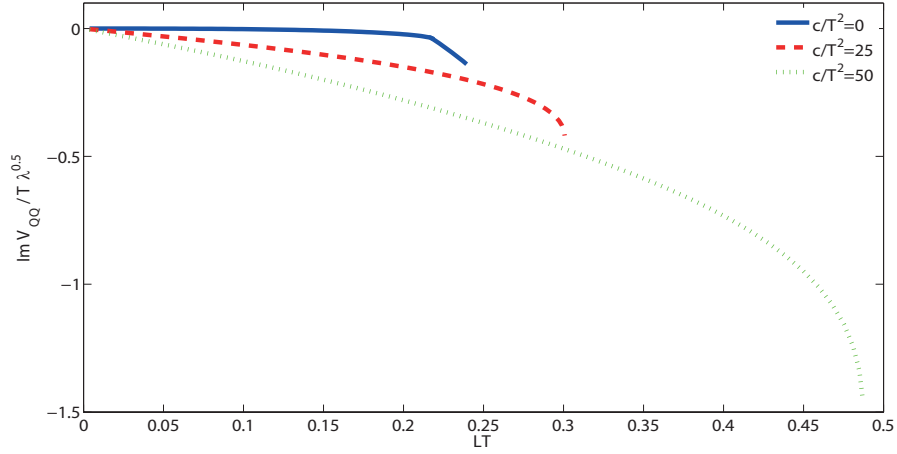


Figure 7. Imaginary part of potential as a function of LT , at a fixed velocity $\eta = 0.4$. The pair is oriented to the hot wind and different values of scaled deformation parameter are contributed. The solid blue curve corresponds to $\frac{c}{T^2} = 0$, the dashed red curve to $\frac{c}{T^2} = 25$ and the dotted green curve to $\frac{c}{T^2} = 50$.

There are two constants of motion which are,

$$\mathcal{H} \equiv Q \equiv \mathcal{L} - \frac{dz}{d\sigma} \frac{\partial \mathcal{L}}{\partial z'} - \frac{dX_d}{d\sigma} \frac{\partial \mathcal{L}}{\partial X'_d} \quad (3.4)$$

$$K \equiv \frac{\partial \mathcal{L}}{\partial X'_d}, \quad (3.5)$$

with (3.3), (3.4) and (3.5) after some algebra one can arrive at,

$$Q^2 [M(z) \cosh^2 \eta - N(z) \sinh^2 \eta] z'(\sigma)^2 + Q^2 V(z) X'_{d-1}(\sigma)^2 + [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta] \{Q^2 - [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta]\} = 0 \quad (3.6)$$

$$K^2 [M(z) \cosh^2 \eta - N(z) \sinh^2 \eta] z'(\sigma)^2 + V(z)(K^2 - V(z)) X_{d-1}'^2(\sigma) + K^2 [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta] = 0. \quad (3.7)$$

with inserting $X_d'^2$ from (3.6) into (3.7) and doing some manipulations the result is,

$$Q^2 V(z) [M(z) \cosh^2 \eta - N(z) \sinh^2 \eta] z'(\sigma)^2 = (V(z) - K^2) [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta]^2 - V(z) [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta] Q^2, \quad (3.8)$$

and

$$Q^2 V^2 (X'_{d-1})^2 = K^2 [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta]^2. \quad (3.9)$$

It is clear that we must have $Z(\sigma = 0) = Z_*$, $Z'(\sigma = 0) = 0$ and $X_d(\sigma = 0) = 0$ so,

$$(V_* - K^2)(V_* \cosh^2 \eta - P_* \sinh^2 \eta) - V_* Q^2 = 0. \quad (3.10)$$

Proceeding by boundary conditions (3.1) and equations of motion (3.8) and (3.9) we arrive at these two relations,

$$\begin{aligned} \frac{L}{2} \sin \theta = & -Q \int_0^{z_*} dz \left\{ \frac{V(z)}{V(z) \cosh^2 \eta - P(z) \sinh^2 \eta} \times \right. \\ & \left. \times \frac{M(z) \cosh^2 \eta - N(z) \sinh^2 \eta}{[(V(z) - K^2) [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta] - V(z) Q^2]} \right\}^{-1/2}, \end{aligned} \quad (3.11)$$

$$\frac{L}{2} \cos \theta = -K \int_0^{z_*} dz \sqrt{\frac{[M(z) \cosh^2 \eta - N(z) \sinh^2 \eta] [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta]}{V(z) \{(V(z) - K^2) [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta] - V(z) Q^2\}}}. \quad (3.12)$$

Finally the action is,

$$S = -\frac{\mathcal{T}}{\pi \alpha'} \int_0^{z_*} dz \sqrt{\frac{V(z) [M(z) \cosh^2 \eta - N(z) \sinh^2 \eta] [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta]}{\{(V(z) - K^2) [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta] - V(z) Q^2\}}}. \quad (3.13)$$

After regularizing it we have,

$$\begin{aligned} S_{\text{reg}} = & -\frac{\mathcal{T}}{\pi \alpha'} \int_0^{z_*} dz \left\{ \sqrt{\frac{V(z) [M(z) \cosh^2 \eta - N(z) \sinh^2 \eta] [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta]}{\{(V(z) - K^2) [V(z) \cosh^2 \eta - P(z) \sinh^2 \eta] - V(z) Q^2\}}} \right. \\ & \left. - \sqrt{M_0(z)} \right\} - \frac{\mathcal{T}}{\pi \alpha'} \int_{z_h}^{z_*} dz \sqrt{M_0(z)}. \end{aligned} \quad (3.14)$$

As before, in the absence of black brane $M(z)$ for $T = 0$ leads to M_0 and $ReV_{Q\bar{Q}} = \frac{S_{\text{str}}^{\text{reg}}}{T}$. For imaginary part we have two degrees of freedom $Z(\sigma)$ and $X_{d-1}(\sigma)$. The string partition function is,

$$Z_{\text{str}} \sim \int D(\delta z) D(\delta X_{d-1}) e^{iS_{\text{str}}(\bar{z} + \delta z, \bar{X}_{d-1} + \delta X_d)}, \quad (3.15)$$

where fluctuations $\delta z(\sigma)$ and $\delta X_{d-1}(\sigma)$ are considered with $\frac{\partial z}{\partial \sigma} \rightarrow 0$ and $\frac{\partial X_{d-1}}{\partial \sigma} \rightarrow 0$.

As before with action (3.2) and partitioning the interval in $2N$ subintervals we arrive at,

$$Z_{\text{str}} \sim \left(\int_{-\infty}^{\infty} d(\delta z_{-N}) d(\delta X_{d-1, -N}) \right) \cdots \left(\int_{-\infty}^{\infty} d(\delta z_N) d(\delta X_{d-1, N}) \right) e^{i \frac{T \Delta x}{2\pi \alpha'} \mathcal{L}_j}, \quad (3.16)$$

and

$$\mathcal{L}_j = \sqrt{\tilde{M}(z(x_j))(z'(x_j))^2 + V(z(x_j))(X'_{d-1}(x_j))^2 + \tilde{V}(x_j)}. \quad (3.17)$$

We expand the classical solution $\bar{z}(0)$ around $\sigma = 0$ to quadratic order on σ . If the string did not sag, then we would have $X_{d-1}(\sigma) = \frac{\sigma}{\tan \tilde{\theta}}$ and around $\sigma = 0$ we will have,

$$X_{d-1}(\sigma) = \frac{\sigma}{\tan \tilde{\theta}} + b\sigma^3 + O(\sigma^5), \quad (3.18)$$

$\tilde{\theta}$ is equal to θ and b is a constant. Because of the symmetry of the problem under reflections with respect to the origin of the (X_1, X_d) plane, $X_d(\sigma)$ must be an odd function of σ so,

$$X'_{d-1}(\sigma)^2 = \frac{1}{\tan^2 \tilde{\theta}} + \frac{6b}{\tan \tilde{\theta}} \sigma^2. \quad (3.19)$$

With inserting (3.19) into (3.17), one can arrive at,

$$\mathcal{L}_j = \sqrt{\tilde{C}_1 x_j^2 + \tilde{C}_2}, \quad (3.20)$$

with these definition,

$$\tilde{C}_1 \equiv \tilde{M}_* \bar{z}''(0)^2 + \frac{1}{2} \left(\frac{V'_*}{\tan^2 \tilde{\theta}} + \tilde{V}'_* \right) \bar{z}''(0) + \frac{6b}{\tan \tilde{\theta}} V_* \quad (3.21)$$

$$\tilde{C}_2 \equiv \left(\frac{V_*}{\tan^2 \tilde{\theta}} + \tilde{V}_* \right) + \left(\frac{V'_*}{\tan^2 \tilde{\theta}} + \tilde{V}'_* \right) \delta z + \left(\frac{V''_*}{\tan^2 \tilde{\theta}} + \tilde{V}''_* \right) \frac{(\delta z)^2}{2}. \quad (3.22)$$

As previous section after some algebraic calculations the explicit analytical expression for $ImV_{Q\bar{Q}}$ would be,

$$ImV_{Q\bar{Q}} = -\frac{1}{4\alpha'} \frac{1}{\sqrt{\tilde{C}_1}} \left[\frac{\left(\frac{V'_*}{\tan^2 \tilde{\theta}} + \tilde{V}'_* \right)^2}{2 \left(\frac{V''_*}{\tan^2 \tilde{\theta}} + \tilde{V}''_* \right)} - \left(\frac{V_*}{\tan^2 \tilde{\theta}} + \tilde{V}_* \right) \right]. \quad (3.23)$$

Again we emphasize that all above derivations about imaginary and real parts of potential are presented in references that we mentioned before, but we follow them here for

convenience of the reader. Now we can come back to our main case and follow it with metric (2.1). With using (3.8) and (3.9) we will have,

$$q^2 \left(\frac{dy}{d\tilde{\sigma}} \right)^2 = (y^4 - \cosh^2 \eta) \left(\left(e^{\frac{c}{\pi^2 T^2 y^2}} \right) (y^4 - 1) - p^2 \right) - q^2 (y^4 - 1) \quad \text{and} \quad (3.24)$$

$$\left(\frac{d\chi}{d\tilde{\sigma}} \right)^2 = \frac{p^2}{q^2} \left(\frac{y^4 - \cosh^2 \eta}{y^4 - 1} \right)^2, \quad (3.25)$$

where we defined the dimensionless variables $y \equiv z_h/z$, $\chi \equiv X_d/z_h$ and $\tilde{\sigma} \equiv \sigma/z_h$ as well as the dimensionless integration constants $q^2 \equiv Q^2 z_h^4/R^4$ and $p^2 = K^2 z_h^4/R^4$ also the boundary conditions become,

$$\begin{aligned} y \left(\pm \pi \frac{LT}{2} \sin \theta \right) &= 0 \\ \chi \left(\pm \pi \frac{LT}{2} \sin \theta \right) &= \pm \pi \frac{LT}{2} \cos \theta. \end{aligned} \quad (3.26)$$

So, (3.11), (3.12) and (3.10) lead to (3.27), (3.28) and (3.29),

$$\frac{LT}{2} \pi \sin \theta = q \int_{y_*}^{\tilde{\Lambda}} \frac{dy}{\sqrt{\left(\left(e^{\frac{c}{\pi^2 T^2 y^2}} \right) (y^4 - 1) - p^2 \right) (y^4 - \cosh^2 \eta) - q^2 (y^4 - 1)}} \quad \text{and} \quad (3.27)$$

$$\frac{LT}{2} \pi \cos \theta = p \int_{y_*}^{\tilde{\Lambda}} dy \frac{y^4 - \cosh^2 \eta}{y^4 - 1} \frac{1}{\sqrt{\left(\left(e^{\frac{c}{\pi^2 T^2 y^2}} \right) (y^4 - 1) - p^2 \right) (y^4 - \cosh^2 \eta) - q^2 (y^4 - 1)}}. \quad (3.28)$$

$$\left(\left(e^{\frac{c}{\pi^2 T^2 y_*^2}} \right) (y_*^4 - 1) - p^2 \right) (y_*^4 - \cosh^2 \eta) - q^2 (y_*^4 - 1) = 0. \quad (3.29)$$

Therefore the real part of potential is,

$$\begin{aligned} \frac{\text{Re } V_{Q\bar{Q}}}{T\sqrt{\lambda}} &= \int_{y_*}^{\infty} dy \left[\frac{e^{\frac{c}{\pi^2 T^2 y^2}} (y^4 - \cosh^2 \eta)}{\sqrt{(y^4 - \cosh^2 \eta) \left(e^{\frac{c}{\pi^2 T^2 y^2}} (y^4 - 1) - p^2 \right) - q^2 (y^4 - 1)}} - e^{\frac{c}{2\pi^2 T^2 y^2}} \right] \\ &\quad - \int_0^{y_*} dy e^{\frac{c}{2\pi^2 T^2 y^2}}. \end{aligned} \quad (3.30)$$

And from (3.23) we arrive at imaginary part of potential as,

$$\begin{aligned} \frac{\text{Im } V_{Q\bar{Q}}}{T\sqrt{\lambda}} &= -\frac{\pi}{4} e^{\frac{c}{2\pi^2 T^2 y_*^2}} \\ &\quad \frac{\left[\frac{2cy_*}{\pi T} - 4\pi T y_*^3 - \frac{2c}{\pi T y_*^3} (\cos^2 \tilde{\theta} + \cosh^2 \eta \sin^2 \tilde{\theta}) \right]^2}{2 \left[20y_*^2 \pi^2 T^2 - 14c - \frac{4c^2}{\pi^2 T^2 y_*^2} - \left(\frac{4c^2}{\pi^2 T^2 y_*^6} + \frac{2c}{y_*^4} \right) (\cos^2 \tilde{\theta} + \cosh^2 \eta \sin^2 \tilde{\theta}) \right]} - \left(y_*^4 - (\cos^2 \tilde{\theta} + \cosh^2 \eta \sin^2 \tilde{\theta}) \right) \\ &\quad \times \frac{1}{\sqrt{y''(0)^2 \left(\frac{y_*^4 - \cosh^2 \eta}{y_*^4 - 1} \right) + \left[\frac{2cy_*}{\pi^2 T^2} - 4y_*^3 - \frac{2c}{y_*^3 \pi^2 T^2} (\cos^2 \tilde{\theta} + \cosh^2 \eta \sin^2 \tilde{\theta}) \right] \frac{1}{2 \sin \tilde{\theta}} y''(0) + \frac{6\tilde{b}}{\tan \tilde{\theta}} (y_*^4 - 1)}}. \end{aligned} \quad (3.31)$$

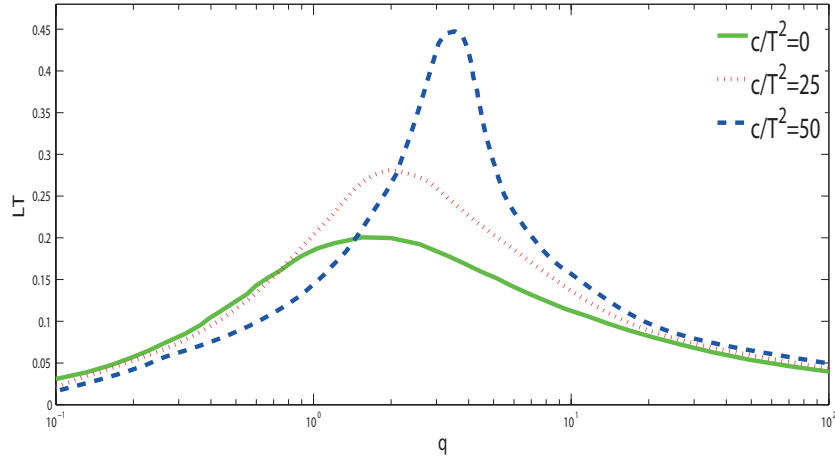


Figure 8. LT as a function of q at $\eta = 1$ and $\theta = \frac{\pi}{3}$. different values of deformation parameter are contributed. The solid green curve corresponds to $\frac{c}{T^2} = 0$, the dotted red curve to $\frac{c}{T^2} = 25$ and the dashed blue curve to $\frac{c}{T^2} = 50$.

We proceed by solving (3.29) numerically to have y_* as a function of q and p , then (3.27), (3.28), (3.30), (3.31) will be functions of p and q . On the other hand for finding p as a function of q , one can solve (3.27) and (3.28) for fixed θ , after doing all these, LT as a function of q is in hand. Before we start to calculate $ImV_{Q\bar{Q}}$ we should obtain $\tilde{\theta}$, $y''(0)$ and \tilde{b} . The $\tilde{\theta}$ is obtained from (3.25) at $\tilde{\sigma} = 0$ and $y = y_*$. We should solve (3.24) and (3.25) with use of boundary conditions (3.26), to evaluate $y''(0)$ and \tilde{b} . After doing all these calculations numerically, with $y''(0)$ and \tilde{b} known, we can calculate $ImV_{Q\bar{Q}}$ as a function of q . So, we will survey $LT(q)$, also real and imaginary parts of potential as a function of LT .

The result of cases with a fixed η and different choices of θ besides a fixed θ and different choices of η have been studeid in [40], so we proceed by fixing both of them and choosing different values of deformation parameter.

In figure 8 we show LT as a function of q for a fixed orientation of the dipole, fixed η and different values of deformation parameter. we know LT_{\max} depends strongly on the rapidity η and it decreases with increasing η [40]. In our plots, we can see that LT_{\max} which indicates the limit of validity of classical gravity calculation, increases with increasing deformation parameter.

In figures 9, 10 and 11 we present $ReV_{Q\bar{Q}}$ as a function of LT . We can see for small values of LT which means short distances or small temperatures, there is a difference between $c = 0$ and $c \neq 0$ cases. As we expected, when deformation parameter contributes to the calculation, the interaction of the pair is relevant with plasma, it is similar to the result of perpendicular case. The other point is that real part of potential has no intense alteration with varying angle for any value of deformation parameter.

In figure 12 we can see $ImV_{Q\bar{Q}}$ as a function of LT . It shows that for angle $\theta < \frac{\pi}{2}$ with decreasing angle, the imaginary part of potential becomes smaller for any values of deformation parameter.

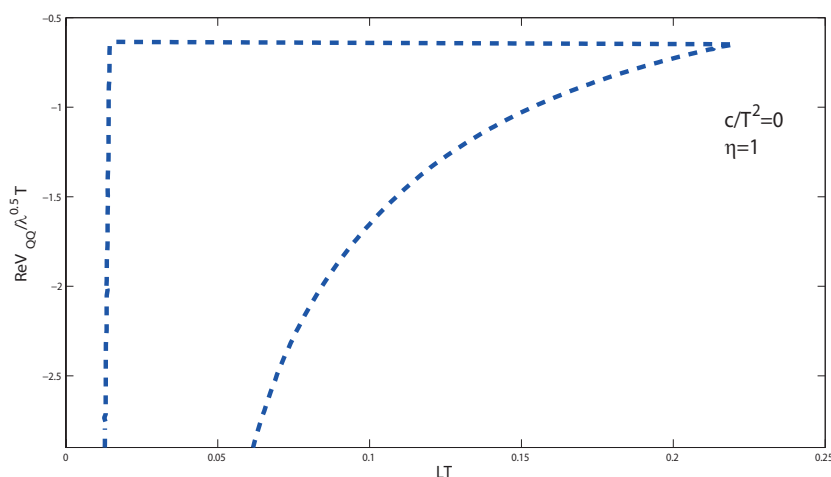


Figure 9. $ReV_{Q\bar{Q}}$ as a function of LT at $\eta = 1$ and $\theta = \frac{\pi}{3}$ and deformation parameter is zero.

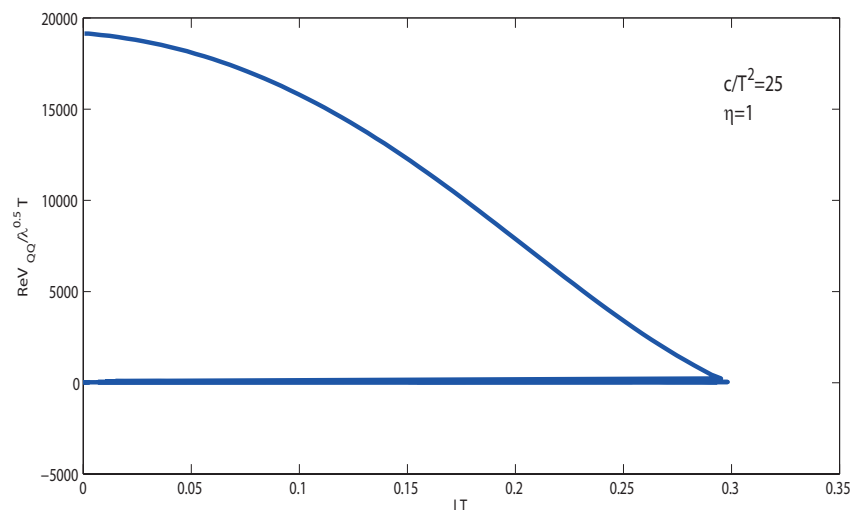


Figure 10. $ReV_{Q\bar{Q}}$ as a function of LT at $\eta = 1$ and $\theta = \frac{\pi}{3}$ and scaled deformation parameter is $\frac{c}{T^2} = 25$.

4 Conclusion

In this article, we have used the method of [40] to investigate the real and imaginary parts of potential for moving heavy quarkonia in plasma with a gravity dual which has deformation parameter in warp factor. At the first step we considered $Q\bar{Q}$ pair oriented perpendicularly to the hot wind and after that we extended all calculations to arbitrary angles. We saw that for both perpendicular and arbitrary angle cases, the limit of classical gravity calculation increases with increasing deformation parameter. Also for nonzero values of c the pair feels moving plasma even in short distances, but for $c = 0$ case the pair does not feel moving plasma at some small values of LT as we expected. We indicated when

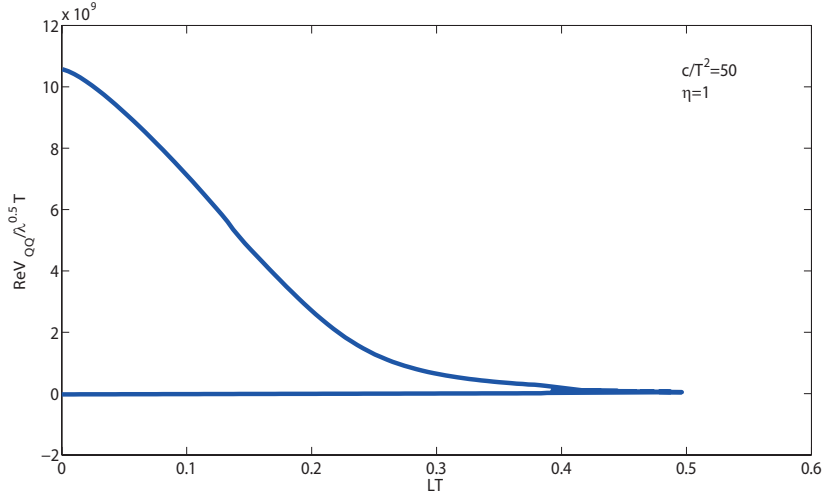


Figure 11. $ReV_{Q\bar{Q}}$ as a function of LT at $\eta = 1$ and $\theta = \frac{\pi}{3}$ and scaled deformation parameter is $\frac{c}{T^2} = 50$.

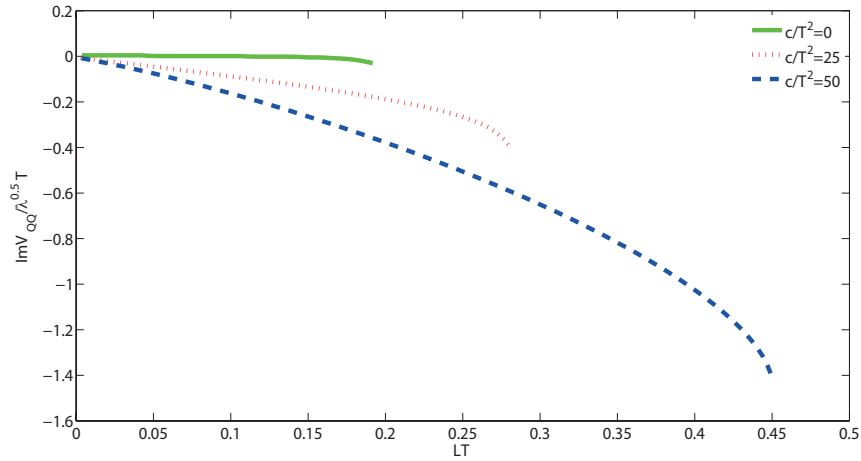


Figure 12. Imaginary part of potential as a function of LT , at a fixed velocity $\eta = 1$ and $\theta = \frac{\pi}{3}$. different values of scaled deformation parameter are contributed. The solid green curve corresponds to $\frac{c}{T^2} = 0$, the dotted red curve to $\frac{c}{T^2} = 25$ and the dashed blue curve to $\frac{c}{T^2} = 50$.

nonzero values of deformation parameter contribute to the imaginary part of potential, the thermal width of quarkonia increases with increasing deformation parameter. Results of perpendicular case in compare with arbitrary angle $\theta < \frac{\pi}{2}$ showed that with decreasing angle, the imaginary part of potential becomes smaller for any values of deformation parameter, but real part of potential has no intense alteration with varying angle for any value of deformation parameter.

Another interesting problem is instead of using the soft wall model we use hyperscaling violation metric background and discuss the moving mesons and investigate real and imaginary parts of potential. This problem with corresponding metric background for the moving meson in plasma media is in hand.

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